

1 The heat equation.

We solve the boundary eigenvalue problem¹:

$$\begin{aligned}\frac{\partial u(x,t)}{\partial t} - k \frac{\partial^2 u(x,t)}{\partial x^2} &= 0, & 0 < x < l, \quad t > 0, \\ u(0,t) = u(l,t) &= 0, & t > 0, \\ u(x,0) &= x,\end{aligned}$$

where k and l are physical constants. By the separation of variables, we assume a solution of the form $u(x,t) = X(x)T(t)$, where $X(x)$ represents the spatial and $T(t)$ the temporal component, respectively. Upon substitution in the equation we have

$$\frac{\dot{T}(t)}{T(t)} = k \frac{X''(x)}{X(x)} = -\lambda,$$

where the functions are equal only if both are equal to some constant denoted by $-\lambda$. The minus sign is for convenience as we seek an imaginary solution that is convertible to a series of sines and cosines. Generally one chooses the spatial ordinary differential equation first because its eigenfunctions will be more fundamental in the resulting separation of variables solution.

Rearranging in a more familiar form, the ODE is

$$X''(x) = \frac{-\lambda}{k} X(x).$$

Letting $D = \frac{d}{dx}$ and substituting back into the equation, we have

$$D^2 = \frac{-\lambda}{k},$$

or

$$D = \pm \sqrt{\frac{-\lambda}{k}} = \pm i \sqrt{\frac{\lambda}{k}}.$$

Then

$$\begin{aligned}X(x) &= c'_1 e^{i\sqrt{\frac{\lambda}{k}}x} + c'_2 e^{-i\sqrt{\frac{\lambda}{k}}x} \\ &= c'_1 \left(\cos \sqrt{\frac{\lambda}{k}}x + i \sin \sqrt{\frac{\lambda}{k}}x \right) + c'_2 \left(\cos \sqrt{\frac{\lambda}{k}}x - i \sin \sqrt{\frac{\lambda}{k}}x \right)\end{aligned}$$

where c'_1 and c'_2 denote arbitrary imaginary constants. Multiplying through and combining, the solution is

$$X(x) = c_1 \cos \sqrt{\frac{\lambda}{k}}x + c_2 \sin \sqrt{\frac{\lambda}{k}}x,$$

where $c_1 = c'_1 + c'_2$ and $c_2 = i(c'_1 - c'_2)$.

¹See Gustafson pp. 38-9.

Applying the boundary conditions at $u(0)$ and $u(l)$ we have

$$0 = c_1$$

and

$$0 = c_2 \sin \sqrt{\frac{\lambda}{k}} l$$

or

$$\lambda_n = \frac{n^2 \pi^2}{l^2} k,$$

where the eigenvalues correspond to integers $n = 1, 2, 3, \dots$, and where the arcsine of zero is equal to $\pm n\pi$. These acquired eigenvalues λ_n determine the temporal factors which we immediately integrate to get:

$$\frac{dT(t)}{T(t)} = -\lambda T(t),$$

hence

$$T_n(t) = d_n e^{-\lambda_n t}.$$

Substituting the results:

$$u(x, t) = \sum_{n=1}^{\infty} c_n \sin \frac{n\pi x}{l} e^{-\frac{n^2 \pi^2 k}{l^2} t}, \quad (1)$$

where the constant d_n was absorbed into c_n . Forming the linear superposition provides our formal separation of variables solution. Applying the third boundary condition ($u(x, 0)$), we have

$$x = \sum_{n=1}^{\infty} c_n \sin \frac{n\pi x}{l} e^{-\frac{n^2 \pi^2 k}{l^2} t}.$$

Solving for the Fourier coefficients we multiply both sides by $\sin \frac{m\pi x}{l}$ and integrate (because if the functions are equal, the integral of the functions are equal too) to get

$$\int_0^l x \sin \frac{m\pi x}{l} dx = \int_0^l c_n \sin \frac{n\pi x}{l} \sin \frac{m\pi x}{l} dx.$$

Recall the right-hand side (RHS) integral is zero if $m \neq n$ and not zero if $m = n$. Letting $m = n$ we have

$$\int_0^l x \sin \frac{n\pi x}{l} dx = \int_0^l c_n \sin^2 \frac{n\pi x}{l} dx.$$

A simple way of evaluating the RHS is to observe that $\frac{\sin^2 \theta + \cos^2 \theta}{2} = 1$. Adding $\cos^2 \frac{n\pi x}{l}$ and dividing by 2, the RHS integral reduces to

$$c_n \int_0^l \frac{1}{2} dx.$$

Integrating the RHS and solving for the Fourier coefficient

$$c_n = \frac{2}{l} \int_0^l x \sin \frac{n\pi x}{l} dx.$$

Integrating by parts we let $u = x$, $du = dx$, $dv = \sin \frac{n\pi x}{l}$, and $v = -\frac{l}{n\pi} \cos \frac{n\pi x}{l}$. Then

$$\begin{aligned} c_n &= \frac{2}{l} \left[-\frac{xl}{n\pi} \cos \frac{n\pi x}{l} + \frac{l^2}{n^2\pi^2} \sin \frac{n\pi x}{l} \right] \Big|_0^l \\ &= \frac{-2l}{n\pi} \cos n\pi = \frac{2l}{n\pi} (-1)^{n+1}. \end{aligned}$$

The final solution is then

$$u(x, t) = \frac{2l}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} \sin \frac{n\pi x}{l} e^{-\frac{n^2\pi^2 k}{l^2} t}. \quad (2)$$

Note that the physical scales are reflected in the eigenfunctions, decay rates, and Fourier coefficients of the solution.